

Interpolation of Markov maps on quantum Orlicz spaces

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OAQD – 12-14 July 2017

Outline

- 1 Classical Orlicz spaces
- 2 Mildly noncommutative function spaces
- 3 Applications
- 4 Wildly noncommutative spaces
- 5 Quantum Orlicz dynamics
- 6 Bibliography

Orlicz functions

Orlicz function: A convex function $\psi : [0, \infty) \rightarrow [0, \infty]$ satisfying

- $\psi(0) = 0$ and $\lim_{u \rightarrow \infty} \psi(u) = \infty$,
- neither identically zero nor infinite valued on all of $(0, \infty)$,
- left continuous at $b_\psi = \sup\{u > 0 : \psi(u) < \infty\}$.

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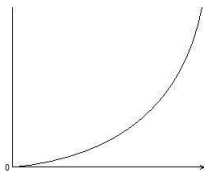
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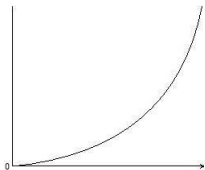


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Defining Orlicz spaces

L^0 the space of all measurable functions on σ -finite (Ω, Σ, m) .

Definition (Orlicz space corresponding to ψ)

$f \in L^0$ belongs to $L^\psi \Leftrightarrow \psi(\lambda|f|)$ is integrable for some $\lambda = \lambda(f) > 0$.

Luxemburg-Nakano norm: $\|f\|_\psi = \inf\{\lambda > 0 : \|\psi(|f|/\lambda)\|_1 \leq 1\}$.

Orlicz norm: $\|f\|_\psi^O = \sup\{|\int_\Omega fg \, dm| : g \in L^{\psi^*}, \|g\|_{\psi^*} \leq 1\}$.

Notational convention: L^ψ (Luxemburg norm); L_ψ (Orlicz norm).

Köthe duality: A measurable function f belongs to $L_{\psi^*}(X, \Sigma, \nu)$ if and only if $fg \in L^1$ for every $g \in L^\psi$.

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New norms on $M_n(\mathbb{C})$

Example: Replace L^∞ by $M_n(\mathbb{C})$, and $\int \cdot d\nu$ by Tr , and see what happens:

- $L^p(M_n(\mathbb{C}), \text{Tr})$ is just $M_n(\mathbb{C})$ equipped with the norm $\text{Tr}(|a|^p)^{1/p}$.
- Similarly $L^\psi(M_n(\mathbb{C}), \text{Tr})$ is $M_n(\mathbb{C})$ equipped with the norm $\|a\|_\psi = \inf\{\lambda > 0 : \text{Tr}(\psi(|a|/\lambda)) \leq 1\}$.

J von Neumann, Some matrix inequalities and metrization of matrix space, *Tomsk Univ Rev* 1(1937), 286-300

Upping the ante: If we play essentially the same game but using a semifinite von Neumann algebra M and an associated fns trace τ instead of $(M_n(\mathbb{C}), \text{Tr})$, the theory still works.



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- M a von Neumann algebra, equipped with a faithful normal semifinite trace $\tau_M = \tau : M^+ \rightarrow [0, \infty]$.
- \tilde{M} the algebra of τ_M -measurable operators: operators affiliated to M , such that for every $\varepsilon > 0$ there exists a projection $e \in M$ with $\tau(e) \leq \varepsilon$, and $a(\mathbb{1} - e) \in M$.

\tilde{M} plays the role of the completion of L^∞ under the topology of convergence in measure.

$f \in \tilde{M}$ belongs to $L^p(M, \tau) \Leftrightarrow \tau(|f|^p) < \infty$ with $\|f\|_p = \tau(|f|^p)^{1/p}$.

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Let f be a fixed element (state) in
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Definition (Pistone-Sempi, 1995)

A measurable function u is said to be a **regular observable** (with respect to f) if the function $\hat{u}(t) = \int e^{tu} f.d\nu$ exists in a neighbourhood of 0, and $\int u f d\nu = 0$.

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Proposition (L. Majewski, 2014)

Let M be a semifinite algebra and $f \in L^1 \cap L \log(L+1)(M, \tau)$ with $f \geq 0$. Then $\tau(f \log(f + \epsilon))$ is well defined for any $\epsilon > 0$.

Moreover

$$\tau(f \log f)$$

is bounded above, and if in addition $f \in L^{1/2}$, it is also bounded from below.

Here $L \log(L+1)(M, \tau)$ is the Orlicz space corresponding to the function $\Psi(t) = t \log(t+1)$.

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Inspired by the controversial work of Boltzmann on the dynamics of rarefied gases [1872], von Neumann expressed entropy as $\text{Tr}(\rho \log(\rho))$ in the context of $B(H)$ (here ρ is a norm 1 element of $\mathcal{S}^1(H)^+$ representing the state of the system).

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So why does the space $L \log(L + 1)(M, \tau) \cap L^1(M, \tau)$ not feature in the context of the pair $\langle \mathcal{S}^1(H), B(H) \rangle$?

In the case of $M = B(H)$, $\tau = \text{Tr}$, we have that

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Implications

Achieved results: $L^{\cosh^{-1}}(M, \tau)$ is a home for regular quantum observables, and $L \log(L + 1)(M, \tau) \cap L^1(M, \tau)$ a home for states with good entropy.

Deeper truths: The space $L^{\cosh^{-1}}(M, \tau)$ is actually an isomorphic copy of the Banach space dual of $L \log(L + 1)(M, \tau)$. So up to isomorphism, $\langle L \log(L + 1)(M, \tau), L^{\cosh^{-1}}(M, \tau) \rangle$ is a dual pair.

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The strange ways of type III L^p spaces

Commutative

$$A = L^\infty(X, \Sigma, \nu) \otimes L^\infty(\mathbb{R})$$

$$\theta_s(f \otimes g)(x, t) = f(x)g(t - s)$$

$$\int \cdot d\nu \otimes \int_{\mathbb{R}} \cdot e^{-t} dt$$

Quantum

“enlarge” M by passing to
 $A = M \rtimes_{\nu} \mathbb{R}$

a dual action of \mathbb{R} on A in the
 form of a group of $*$ -auto-
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(Haagerup, 1979): For any measurable function f on X (finite ν -almost everywhere) we have that

$$f \otimes e^{(\cdot)/p} \in \tilde{A} \iff f \in L^p(X, \Sigma, \nu).$$



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(Haagerup, 1979): By analogy with the classical setting, we may define

$$L^p(M) = \{a \in \tilde{A} : \theta_s(a) = e^{-s/p} a \text{ for all } s \in \mathbb{R}\}.$$



Constructing $A = M \rtimes_{\nu} \mathbb{R}$

- Replace \mathcal{H} with $L^2(\mathbb{R}, \mathcal{H})$.
- The map $a \rightarrow \pi(a)$ defines an embedding of M into $B(L^2(\mathbb{R}, \mathcal{H}))$, where $(\pi(a)(\eta))(t) = \sigma_{-t}^{\nu}(a)(\eta(t))$ for all $a \in M$ and all $\eta \in L^2(\mathbb{R}, \mathcal{H})$.
- Throw in some shift operators $(\lambda(s)(\eta))(t) = \eta(t - s)$,
- and generate the von Neumann algebra $A = M \rtimes_{\nu} \mathbb{R} \subset B(L^2(\mathbb{R}, \mathcal{H}))$ from these two classes of maps.

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- The map $a \rightarrow \pi(a)$ defines an embedding of M into $B(L^2(\mathbb{R}, \mathcal{H}))$, where $(\pi(a)(\eta))(t) = \sigma_{-t}^{\nu}(a)(\eta(t))$ for all $a \in M$ and all $\eta \in L^2(\mathbb{R}, \mathcal{H})$.
- Throw in some shift operators $(\lambda(s)(\eta))(t) = \eta(t - s)$,
- and generate the von Neumann algebra $A = M \rtimes_{\nu} \mathbb{R} \subset B(L^2(\mathbb{R}, \mathcal{H}))$ from these two classes of maps.

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Type III Orlicz spaces

Haagerup's construction of L^p -spaces for type III von Neumann algebras can be extended to also allow for the construction of Orlicz spaces. (L, 2014)

The classical roots of the construction: Let $M = L^\infty(X, \Sigma, \nu)$, and let $A = L^\infty(X, \Sigma, \nu) \otimes L^\infty(\mathbb{R})$ be as before.

Given an Orlicz function Ψ , define $\varphi_\Psi : [0, \infty) \rightarrow [0, \infty)$ by

$$\varphi_\Psi(t) = \frac{1}{\Psi^{-1}(1/t)}.$$

For any measurable function f on X , we then have that

$$f \otimes \varphi_\Psi(e^{(\cdot)}) \in \tilde{A} \iff f \in L^\Psi(X, \Sigma, \nu).$$



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Given a von Neumann algebra M with fns weight ν ,

- let $\hat{\nu}$ be the dual weight on the crossed product $A = M \rtimes_{\nu} \mathbb{R}$,
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Emergent challenge

Challenge: Given a Markov map T with a canonical action on M and $L^1(M)$, can we show that it has a nice action on a large enough class of Orlicz spaces? First pause to see what is known.

Theorem (Yeaton 1977; HJX 2010)

Let $T : M \rightarrow M$ be a positive map for which there exists some $C_1 > 0$ such that $\nu(T(x)) \leq C_1 \nu(x)$ for all $x \in M^+$. Then for each $1 \leq p < \infty$, T canonically extends to a positive bounded map $T_p : L^p(M) \rightarrow L^p(M)$ such that $\|T_p\| \leq C_\infty^{1-(1/p)} \cdot C_1^{1/p}$ where $C_\infty = \|T(\mathbb{1})\|_\infty$.

Problem: The proof uses complex interpolation. To date complex interpolation does not work for quantum Orlicz spaces.

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An alternative strategy

Challenge 2: Can we overcome the lack of access to complex interpolation, by passing to a smaller class of Markov maps, namely the CP Markov map? If so how?

Idea:

- Show that under acceptable assumptions, $T : M \rightarrow M$ extends to a map \tilde{T} on $A = M \rtimes_{\sigma} \mathbb{R}$,
- and from there to a map on $(L^{\infty} + L^1)(A, \tau_A)$.
- Then see if any of the Orlicz spaces $L^{\psi}(M)$ live inside $(L^{\infty} + L^1)(A, \tau_A)$, and try to extract the action from that.

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HJX extension result

Theorem

Let $T : M \rightarrow M$ be a completely bounded normal map such that $T \circ \sigma_t^\alpha = \sigma_t^\alpha \circ T$, $t \in \mathbb{R}$. Then the prescription

$\tilde{T}(\lambda(s)\pi(x)) = \lambda(s)\pi(T(x))$ ($x \in M, s \in \mathbb{R}$) generates a unique bounded normal extension \tilde{T} of T to $A = M \rtimes_{\sigma^\alpha} \mathbb{R}$ with

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(1) $\tilde{T}(\pi(x)\beta) = \pi(T(x))\beta$ for all $x \in M$ where β is the canonical semifinite normal weight on A generated by $\pi(N)$, $s \in \mathbb{R}$.

(2) $\tilde{T} \circ \sigma_t^\beta = \sigma_t^\beta \circ \tilde{T}$ for all $t \in \mathbb{R}$ (β is the dual weight of π).

(3) $\tilde{T} \geq 0 \Leftrightarrow T \geq 0$.

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- 2 $\tilde{T} \circ \sigma_t^{\hat{\nu}} = \sigma_t^{\hat{\nu}} \circ \tilde{T}$ for all $t \in \mathbb{R}$ ($\hat{\nu}$ is the dual weight of ν).
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$\tilde{T}(\lambda(s)\pi(x)) = \lambda(s)\pi(T(x))$ ($x \in M, s \in \mathbb{R}$) generates a unique bounded normal extension \tilde{T} of T to $A = M \rtimes_{\sigma^\nu} \mathbb{R}$ with

$\|T\| = \|\tilde{T}\|$. Moreover:

- 1 $\tilde{T}(a\pi(x)b) = a\pi(T(x))b$ for all $a, b \in B$ where B is the von Neumann subalgebra generated by all $\lambda(s), s \in \mathbb{R}$.
- 2 $\tilde{T} \circ \sigma_t^{\hat{\nu}} = \sigma_t^{\hat{\nu}} \circ \tilde{T}$ for all $t \in \mathbb{R}$ ($\hat{\nu}$ is the dual weight of ν).
- 3 $T \geq 0 \Rightarrow \tilde{T} \geq 0$.
- 4 $\nu \circ T \leq \nu \Rightarrow \hat{\nu} \circ \tilde{T} \leq \hat{\nu}$.



HJX extension result

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Real interpolation to the rescue

Corollary

Let T and \tilde{T} be as before. If each of (1)-(4) holds, then $\tau_A \circ \tilde{T} \leq \tau_A$ where τ_A is the canonical trace on $A = M \rtimes_{\sigma^v} \mathbb{R}$. The map \tilde{T} then also canonically induces a map on the space $(L^\infty + L^1)(A, \tau_A)$.

Proof The first claim follows by applying some ideas from Pedersen and Takesaki's seminal paper. For the second claim apply Yeadon's ergodic result to see that \tilde{T} induces a bounded map on $L^1(A, \tau_A)$, and then apply real interpolation to get the conclusion.

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Markov dynamics

Definition

Let φ_ψ be the fundamental function of the space $L^\Psi(0, \infty)$, and let $M_\psi(t) = \sup_{s>0} \frac{\varphi_\psi(st)}{\varphi_\psi(s)}$. We call the quantity

$$\bar{\beta}_{L^\Psi} = \inf_{1 < t} \frac{\log M_\psi(s)}{\log s}$$

the upper fundamental index of $L^\Psi(M)$.

Proposition

If $\bar{\beta}_{L^\Psi} < 1$, then $L^\Psi(M) \subset (L^\infty + L^1)(A, \tau_A)$ (isomorphically). Moreover $L^\Psi(M)$ is an invariant subspace of the extension \tilde{T} of T to $(L^\infty + L^1)(A, \tau_A)$. This class includes $L^{\cosh^{-1}}(M)$!! (The space of regular observables.)

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The emergent picture

Based on the preceding analysis we may propose the following framework as an axiomatic foundation for Quantum Statistical Mechanics:

- Corresponding to each quantum system there is a pair (M, ν) (where M is a von Neumann algebra and ν an associated faithful normal semifinite weight) describing the system.
- The pair of spaces $(L \log(L+1)(M), L^{\cosh^{-1}}(M))$ are respectively homes for good states and good observables of this system.
- There is a Dirichlet form \mathcal{E} (representing an energy potential) describing Markov dynamics on the space $L^{\cosh^{-1}}(M)$ of regular observables.

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